

Modelling of a Dynamic Torque Calibration Device and Determination of Model Parameters

Leonard Klaus, Thomas Bruns, Michael Kobusch

Physikalisch-Technische Bundesanstalt (PTB), Bundesallee 100, 38116 Braunschweig, Germany

ABSTRACT

For the dynamic calibration of torque transducers, a calibration device has been developed. This paper describes the model of the measuring device and methods for the determination of its model parameters. The modelling of the calibration set-up is required for the identification of the corresponding model parameters of the torque transducer under test. These parameters describe the transducer's dynamic behaviour. Measurement methods and devices for the determination of mass moment of inertia and torsional stiffness are explained. This research is part of EMRP JRP IND09 – "Traceable Dynamic Measurement of Mechanical Quantities".

Keywords: Dynamic torque calibration, mass moment of inertia, torsional stiffness, sinusoidal excitation, modelling of mechanical systems.

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Corresponding author: Leonard Klaus, leonard.klaus@ptb.de

1. INTRODUCTION

Dynamic torque applications exist in many fields of industry, e.g. engine test stands or power torque tools. Efficiency measurements of combustion engines and electric drives are carried out by means of rotational speed and dynamic torque measurements in engine test stands. The torque output of these drives may be significantly dynamic with a frequency content up to the kilohertz range. From experience with force transducers, it is known that transducers may have a dynamic behaviour which can influence the measuring results.

2. MEASURING DEVICE

At present, static torque calibrations can be carried out in a wide range and with high precision, but there are no facilities and standards to determine the dynamic behaviour of torque transducers. For this reason, a method for primary dynamic torque calibration was developed and a proof-of-principle calibration device was designed and manufactured. This device enables dynamic torque excitations in a frequency range of up to 1 kHz with up to 20 N·m sinusoidal torque.

The measurement principle is in analogy to the periodic force excitation based on Newton's second law [1]

$$M(t) = \ddot{\varphi}(t) \cdot J. \tag{1}$$

Figure 1 shows the components of the calibration device and the corresponding mechanical spring-mass-damper model. The transducer under test is mounted via coupling elements to a rotational exciter. The coupling elements are composed of an upper and a lower half which are connected by a steel diaphragm to reduce parasitic bending moments acting on the transducer. At the same time, they offer a high torsional stiffness. Interchangeable collets enable the mounting of transducers with shaft ends of various diameters in the measuring device. At the top, the transducer carries a known mass moment of inertia (MMOI) I, which is composed of components for an angular acceleration measurement set-up and a radial air bearing. The air bearing reduces the axial load as well as bending moments acting on the transducer under test in conjunction with low friction operation. The primary measurement of the angular acceleration is carried out by means of a radial grating disk and a laser-Doppler interferometer [1].

3. MODEL

The measurement principle of the vast majority of torque transducers is based on strain gauges. This principle is used for force transducers as well, so the implementation of a model for torque transducers was realised based on the experience with force transducers [2]. Strain gauge transducers are designed to measure strain at one or more positions which have a significantly lower stiffness than the rest of the transducer. This part can be described as torsional spring (c_T) for torque transducers. The remaining components of the transducer can be assumed to be rigid and will be allocated to the mass moment of inertia element above (J_H) or below (J_B) the torsional spring, respectively. In parallel to the torsional spring, the model assumes a damping element. The behaviour of the system is supposed to be linear and time invariant (LTI).

The distribution of the head and base MMOI of the transducer under test ($J_{\rm H}$ and $J_{\rm B}$) is dependent on the mechanical design of the transducer.

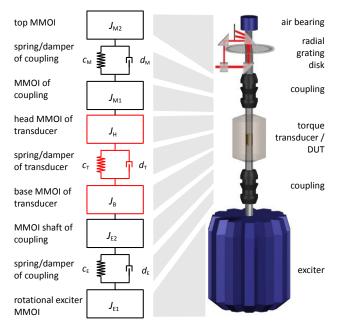


Figure 1. Measuring device and its model.

Because of the fact that torque transducers are always coupled to the mechanical environment at both ends, different coupled MMOI of different stiffness and damping may have influence on the frequency response of the system, which affects the frequency-dependent output signal of the transducer. To be able to identify the model parameters of the transducer, it is necessary to include the entire measurement set-up in the model.

The measuring device is coupled to both ends of the transducer, as is its model. It is assumed to be linear and time invariant, as is the transducer's model.

As the coupling elements cannot be assumed to be totally rigid to torsional loads, these are represented in the model of the measuring device by two torsional spring and damper elements which are coupled to mass moments of inertia. Rigid connections are divided into properties of the transducer and properties of the measuring device (e.g. $J_{\rm E2}$

and $J_{\rm B}$ in Figure 1) in order to be able to determine solely the model parameters of the transducer.

The mass moment of inertia $J_{\rm M2}$ includes all components of the measuring device on top of the upper half of the coupling. $J_{\rm M1}$ includes the lower part of the coupling, as well as adapters for the mounting of the transducer; the same applies to $J_{\rm E2}$ but in the opposite direction.

Only if the dynamic properties of the measuring device are known, will the determination of the transducer's model parameters from measurement data possible. Thus, measurement techniques for the determination of these parameters have to be developed.

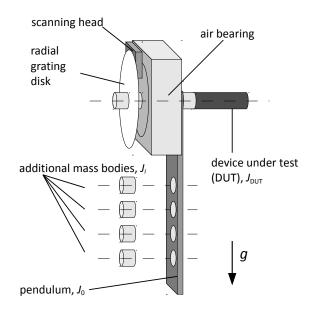


Figure 2. Pendulum for the measurement of the mass moment of inertia.

4. MASS MOMENT OF INERTIA

The mass moment of inertia that generates the dynamic torque acting on the transducer, when applied to an angular acceleration, needs to be determined with high precision.

For this purpose, a dedicated measuring device for the mass moment of inertia has been developed (see Figure 2). Its measurement method is based on the principle of a physical pendulum. For small angles of pendulum excitation, the nonlinearity of the pendulum can be neglected – though this effect has to be taken into account for the measurement uncertainty evaluation. The frequency of the pendulum's swing ω depends on its mass moment of inertia J and the restoring torque consisting of the mass of the pendulum m, the distance of the centre of mass from the axis of rotation l and the gravitational acceleration g

$$J = \frac{m \cdot g \cdot l}{\omega^2}.$$
 (2)

To minimise frictional losses, the pendulum set-up uses an air bearing [3]. The device under test (DUT) is mounted on the axis of rotation. The mass moment of inertia of the DUT is derived from the change of the pendulum's swing frequency. For the measurement of the angular excitation, an optical angle encoder is used. Its grating disk has an angular

pitch of 0.04°. This angular resolution can be additionally increased by interpolation methods.

The measured mass moment of inertia J_{total} is composed of the pendulum's MMOI and the MMOI of all mounted components (e.g. the device under test)

$$J_{\text{total}} = J_{\text{DUT}} + J_0 + \sum J_i.$$
 (3)

As the pendulum's mass moment of inertia J_0 is unknown, it has to be determined first. This is to be carried out in the same procedure as used subsequently for the mounted device under test. The unknown property can be identified by adding known mass moments of inertia to the pendulum. This is to be carried out by mounting auxiliary mass bodies at different distances from the axis of rotation into the pendulum lever.



Figure 3. Measurement set-up for the determination of mass moment of inertia with 3 mounted mass bodies and mounted DUT.

The mass bodies of the pendulum have well-known dimensions and weights. With this information, the mass moment of inertia for rotation around the centre of gravity of each mass body can be calculated. The distance of the mounting position of the masses to the axis of rotation of the pendulum was measured in advance. The acting mass moment of inertia J_i of each mass body at a certain mounting position is given by the Huygens-Steiner theorem.

The mounted mass bodies do not only influence the mass moment of inertia of the pendulum, but also the mass and the centre of mass of the pendulum, i.e. the restoring torque on the right-hand side of equation 2. Each configuration is described by

$$J_0 + \sum J_i = \frac{(m_0 \cdot l_0 \cdot g) + \sum m_i \cdot l_i \cdot g}{4 \cdot \pi^2 \cdot f^2}.$$
 (4)

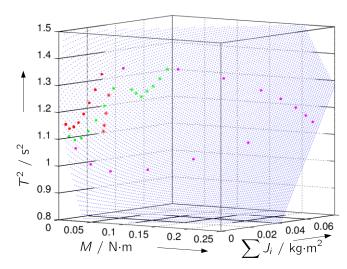


Figure 4. Measurement values for different auxiliary mass bodies mounted at different positions in the pendulum (dots) and result of the bivariate regression (dotted surface).

To determine the swing frequency of the pendulum in a given test configuration, the angle encoder output is acquired for a given time with a sufficiently high sampling rate for a similar range of pendulum excitation angles. After the predetermination of frequency, magnitude, phase and damping, a four parameter Levenberg-Marquardt fit is calculated on a damped sine function

$$y(t) = C \cdot e^{-\delta t} \cdot \sin(\omega t - \varphi), \tag{5}$$

where the predetermined values are used as initial parameters. Experience gathered after the first measurements shows a stable behaviour of the fit algorithm with the measured data and a very good agreement of measured data and fit.

The unknown mass moment of inertia of the pendulum and, depending on the configuration, of the mounted device under test can be derived by varying the configuration of the pendulum by means of mounting different auxiliary mass bodies. Several measurements are necessary to vary the two parameters *mass moment of inertia* and *restoring torque* sufficiently for the identification purpose.

This leads to a solvable bivariate regression problem in the form of

$$\mathbf{y} = f(\mathbf{x}, \mathbf{p}). \tag{6}$$

Every measured frequency f of each configuration is part of the vector of observed values \mathbf{y} ; the added auxiliary mass bodies contribute to the matrix of independent values \mathbf{x} , which consist of $\sum J_i$ and $\sum m_i \cdot l_i \cdot g$. The vector of approximated parameters \mathbf{p} consists of J_0 and $m_0 \cdot l_0 \cdot g$.

For the successful identification of the unknown parameters, is it beneficial to maximize the variation of the independent values. This can be achieved by choosing mass bodies and mounting positions in an optimized way. This effect is illustrated in Figure 4, which shows different mass body configurations.

The pendulum is modelled assuming undamped oscillations (see equation (2)). This simplification yields

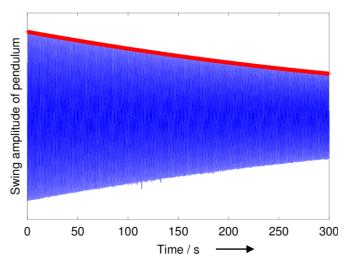


Figure 5. Decay of the pendulum's swing amplitude.

small deviations only for systems with low damping. To determine the damping, the decay of the pendulum's amplitude was measured.

The oscillations of the pendulum were measured by means of a laser-Doppler interferometer, as the angle measurement set-up had not yet been installed at the time of measurement. The interferometer traced the pendulum's swing at the bottom end of the pendulum.

The damping of the pendulum was determined by calculating a non-linear regression approximating the measurement values with a damped sine fit function (see equation (5)). The decay of the pendulum swing amplitude for the measurement data shown in Figure 3 is calculated and equals $\delta = 2.23 \cdot 10^{-3} \, \text{s}^{-1}$. The damping slightly influences the resonant frequency of the pendulum [4]. The higher the damping, the lower the frequency of the pendulum ω_1 in comparison to the frequency of the undamped pendulum ω_0 . This effect has to be taken into account

$$\omega_0 = \sqrt{(\omega_1^2 + \delta^2)} \,. \tag{7}$$

For the described measurement set-up, this effect is in the range of $8 \cdot 10^{-8}$.

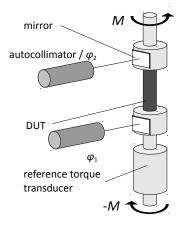


Figure 6. Torsional stiffness measurement set-up.

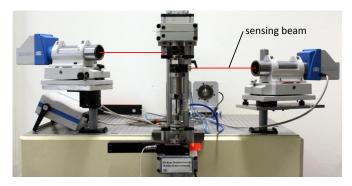


Figure 7. Measurement set-up for torsional stiffness, indicated sensing beams of autocollimators (red).

5. TORSIONAL STIFFNESS

The stiffness of the components of the measuring device has great influence on the torque measurement as it affects the dynamic behaviour of the whole system. The torsional stiffness c is defined as the torque-to-torsion ratio (with torque M and the torsion angle $\Delta \varphi = \varphi_2 - \varphi_1$)

$$c = \frac{M}{\varphi_2 - \varphi_1}. (8)$$

The determination of the torsional stiffness is realised by a measurement set-up (see Figure 6), in which a well-known torque is applied to the DUT while the angle of torsion is measured with high precision by means of two autocollimators.

Two mirrors attached with support clamps to both sides of the DUT provide the sensing surface of the autocollimators. The torque-dependent difference of the two angle readings of both autocollimators yields the torsional angle. The change in the angle of the mirror surface is equal to the torsion. Mirrors with a planarity of $1/10~\lambda$ enable the determination of the angle with uncertainties of less than one arcsecond; the measurement range amounts to rotation angles of up to \pm 600 arcseconds. In this measurement set-up, the statically applied torque is measured by means of a reference torque transducer.

To measure the torsional stiffness, PTB's 20 N·m Torque Calibration Machine was equipped with two electronic autocollimators in order to perform measurements of the torsional angle. Figure 7 shows a photograph of the set-up. The Torque Calibration Machine is seen in the centre, on the right is the autocollimator tracing the bottom mirror, on the left, the autocollimator sensing at the top mirror.

A measurement routine was developed for the evaluation of the torsional stiffness. The applied load steps are based on the DIN 51309 standard for torque transducer calibrations. After preloading to avoid hysteresis behaviour in the torque calibration system, the applied torque (load) increases in steps of 10% to the full load. The number of load steps shown in Figure 8 was increased in comparison to the standard. For the measurement, both a clockwise and a counter-clockwise torsional loading cycle were applied.

The voltage output of the bridge amplifier connected to the reference torque transducer and the angle values of the two autocollimators were acquired simultaneously. Several

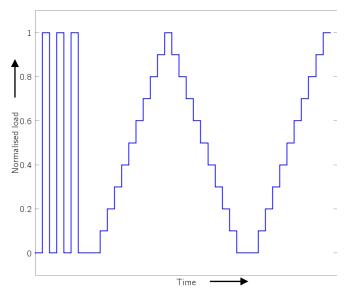


Figure 8. Load steps of the torsional stiffness measurement routine.

hundred measurement values were recorded for each torque step.

The first measurement results are very encouraging, as seen in Figure 9. The torsional angle values for the different torque steps show a linear dependency. The first order regression line excellently fits the mean values of the different torque levels. The value for torsional stiffness results from the gradient of the regression line. The angle offset is given by the initial absolute angle values of the autocollimators.

6. CONCLUSIONS

This paper describes the modelling and the determination of the model parameters of the dynamic torque calibration device at PTB. The modelling of the measuring device is a prerequisite to be able to determine the transducer's dynamic properties.

The described methods enable the measurement of the torsional stiffness and the mass moment of inertia. The torsional stiffness is determined by applying a well-known torque and by measuring the torsion angle on both sides of the DUT. The mass moment of inertia is determined by means of the swing frequency of a compound pendulum. It is intended to determine the model parameters of the components of the dynamic torque calibration device with these measurement methods.

For the future model-based calibration of torque transducers, the parameters of the measuring device need to

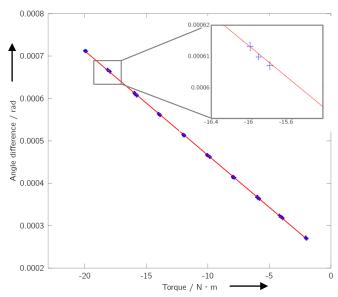


Figure 9. Measurement result with standard deviation (blue) and regression line (red).

be known in order to be able to identify the model parameters of the torque transducer under test from measurement data.

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