Phase correlation functions: FFT vs. FHT

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ABSTRACT
The ability to process an image is a crucial skill in many measurement activities. In image processing or pattern recognition, Fast Fourier Transform (FFT) is widely used. In particular, the Phase Only Correlation (POC) method demonstrates high robustness and subpixel accuracy in pattern matching. However, there is a disadvantage in the required memory machine because of the calculation of 2D-FFT. In applications in which the use of memory is a critical element, Fast Hartley Transform (FHT) seems to be a good substitute. In this context, the use of Hartley’s transform can be of interest for apps implemented on portable systems e.g. smartphones. In this article, we present a comparison of the implementations of the phase correlation function using FFT and FHT. Particular attention is given to the analytical steps necessary to implement the POC by means of the Hartley transform.

1. INTRODUCTION
The ability to accurately process images and video signals is an important skill in many research fields. One of the more widespread techniques used is the Fourier transform [1], [2]. In fact, Fourier’s theorem is not only one of the most impressive results of analysis and one of the most important mathematical tools available to physicists and engineers, but it may be considered an indispensable instrument in digital image processing. In this field, the Fourier’s transform and correlation-based analysis are crucial elements of the pattern-matching process [3], [4]. In many image-processing applications, it is necessary to determine the correspondence between two or more images. The different images can either be images of the same object taken from different sensors or images of the same object taken at different times. This matching analysis is typically performed using either cross-correlation in the spatial domain or phase correlation in the frequency domain. In general, spatial cross-correlation [5], [6] yields several broad peaks and a main peak whose maximum is not always clearly defined; therefore, it is difficult to evaluate and locate it. On the other hand, correlation in the frequency domain is preferable in the case that the images are corrupted by noise.

Figure 1 shows the correlation function between two images, one of which is corrupted by noise.

In this figure, it is possible to note the difference between the cross-correlation in the spatial domain and the phase correlation performed in the frequency domain.

Figure 1. Example of correlation between two images (a) Image ‘1’; (b) Image ‘2’; (c) correlation in the spatial domain; (d) correlation in frequency domain.
The phase correlation in the frequency domain is efficiently performed by Phase-Only Correlation (POC) [7], [8]. This correlation method demonstrates high robustness and accuracy in pattern matching and in image registration. However, there is a disadvantage in terms of the required machine memory because of the calculation of 2D-FFT. The Fast Hartley Transform (FHT) can be a valid alternative to the Fast Fourier Transform (FFT) [9]. The Hartley transform resembles a Fourier transform, but it is free from the need to process complex numbers. The Hartley transform also has better properties and faster algorithms than the Fourier one; therefore, it can represent a valid alternative, particularly useful when a greater efficiency in memory requirements is needed.

In this paper, we analyse the equivalence of the POC computation by means of FFT and FHT. The Hartley transform, in some cases, is a valid alternative to the Fourier transform, in particular in all those applications implemented on embedded systems or smartphones.

The remaining of this paper is organised as follows. The definition of the POC function and its basic properties are described in section 2. The properties of the Hartley transform and the definition of POC in Hartley space are given in section 3. In section 4, we present a set of experiments for evaluating the performance of the proposed methods. Finally, conclusions are drawn in section 5.

2. PHASE-ONLY CORRELATION (POC)

POC is a digital image processing technique that is normally used to estimate the ‘similarity’ between two images. It is extensively used in many practical applications. In particular, POC is employed in image registration [10]-[12], fingerprint matching [13], [14], iris recognition [15], palm recognition [16], PIV analysis [17], digital image stabilisation [18], and security applications [19]-[24].

Phase correlation is based on the well-known Fourier shift property [25]. If we consider two $N_x \times N_y$ digital images $I_x(x,y)$ and $I_y(x,y)$. We assume that the index ranges are $x = -M_x, \ldots, M_x$, and $y = -M_y, \ldots, M_y$ for mathematical simplicity. Hence, $N_x = 2M_x + 1$ and $N_y = 2M_y + 1$. Let $F_x(u,v)$ and $F_y(u,v)$ denote the 2D Discrete Fourier Transforms of the two images $I_x(x,y)$ and $I_y(x,y)$.

$F_x(u,v)$ and $F_y(u,v)$ are given by:

$$F_x(u,v) = \frac{1}{N_x N_y} \sum_{x=-M_x}^{M_x} \sum_{y=-M_y}^{M_y} I_x(x,y) \exp \left[ -j 2\pi \frac{ux + vy}{N_x} \right]$$

(1)

$$F_y(u,v) = \frac{1}{N_x N_y} \sum_{x=-M_x}^{M_x} \sum_{y=-M_y}^{M_y} I_y(x,y) \exp \left[ -j 2\pi \frac{ux + vy}{N_x} \right]$$

(2)

where $u = -M_x, \ldots, M_x$, $v = -M_y, \ldots, M_y$.

The cross-phase spectrum (or normalised cross spectrum) is defined as [26]

$$\frac{F_x(u,v) \cdot F_y^*(u,v)}{|F_x(u,v) \cdot F_y^*(u,v)|} = \exp \left[ -j (ux + vy) \right]$$

(3)

with $F_y^*(u,v)$ being the conjugate function of $F_y(u,v)$. The POC function is the 2D inverse discrete Fourier transform of equation (3). In other words, the POC surface is defined as

$$r(x,y) = F^{-1} \left[ \frac{F_x(u,v) \cdot F_y^*(u,v)}{|F_x(u,v) \cdot F_y^*(u,v)|} \right]$$

(4)

If both images are similar, the POC function $r(x,y)$ gives a distinct sharp peak. Conversely, if the two images are not similar, the peak significantly drops. The height of the peak gives a similarity measurement of the matching, and the position of the peak in the $x \times y$ axis shows the translation displacement of two images. The most remarkable property of POC compared to the ordinary correlation is that it is not influenced by brightness change, and it is highly robust against noise.

POC, in some cases, is defined as [27], [28]

$$r(x,y) = F^{-1} \left[ \frac{F_x(u,v) \cdot F_y^*(u,v)}{|F_x(u,v) \cdot F_y^*(u,v)|} \right]$$

(5)

In equation (5), the coefficient $\alpha$ controls the correlation peak width. Optimum values range from $\alpha = 0$ (for images characterised by high spatial frequency content and a high noise level) to $\alpha = 0.5$ (for low-noise images with a less fine structure). For values greater than 0.5, the high frequency noise is magnified [29].

Figure 2 shows two examples of POC outputs.

In some applications, instead of FFT, it is possible to use FHT. Using FHT, equation (4) and equation (5) can be computed efficiently and without loss of information [30].

3. THE HARTLEY TRANSFORM

The Hartley transform, introduced in 1942 [31], became widely known after the publication of Bracewell’s study in 1983 [32].

It is an integral transform closely related to the Fourier transform; keeps all the useful properties thereof, and can be used to obtain the power spectrum and perform convolution directly through the output real-valued data without firstly calculating the real and imaginary parts. The use of the Hartley transform also reduces the processing time and needs less memory. Therefore, it is particularly well suited for application on mobile devices and embedded systems.

The Hartley transform is one type of integral transformation that is closely related to the Fourier transform. It has the same integral core of the transform and the inverse transform, and it also has most of the characteristics of the Fourier transform.

The Hartley transform can be obtained from the Fourier integral by replacing the exponential function
\[ \exp[-j2\pi(ax+vy)] = \cos[2\pi(ax+vy)] - jsin[2\pi(ax+vy)] \]

with \[ \cos[2\pi(ax+vy)] = \cos(2\pi x) + \sin(2\pi y) \] .

The Hartley transform \( S_h(u,v) \) of a real signal \( s(x,y) \) and its inverse transform are defined as

\[
S_h(u,v) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} s(x,y) \text{cas}[2\pi(ax+vy)] \, dx \, dy
\]

\[
s(x,y) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} S_h(u,v) \text{cas}[2\pi(ax+vy)] \, du \, dv .
\]

The Fourier and Hartley transforms are very similar and share many properties. Consequently, many applications using the Fourier transform can be performed by the Hartley one [33].

To derive the relationship between the Fourier and Hartley transforms, their symmetry must be considered. Let us split the Hartley transform \( S_h(u,v) \) into its even and odd parts \( E_h(u,v) \) and \( O_h(u,v) \), so that

\[
S_h(u,v) = E_h(u,v) + O_h(u,v) .
\]

The even part of the function is what we get by reversing the function (changing \( u \) to \(-u\) and \( v \) to \(-v\)), adding the reversed function to the original and dividing by two, that is

\[
E_h(u,v) = \frac{S_h(u,v) + S_h(-u,-v)}{2} .
\]

The odd part is formed by subtracting the reversed function and dividing by two:

\[
O_h(u,v) = \frac{S_h(u,v) - S_h(-u,-v)}{2} .
\]

Any function may be split uniquely into even and odd parts and from the even and odd parts, if given, the original function may be uniquely reconstituted.

Obviously, the Fourier transform can also be divided into even and odd parts: \( E_f(u,v) \) and \( O_f(u,v) \). Furthermore,

\[
E_f(u,v) = \text{Re}[S_h(u,v)] ,
\]

\[
O_f(u,v) = -\text{Im}[S_h(u,v)] .
\]

By considering the definitions of the Hartley and Fourier transforms, we have \( E_h(u,v) = E_f(u,v) \) and \( O_h(u,v) = -O_f(u,v) \). Therefore,

\[
S_h(u,v) = \text{Re}[S_f(u,v)] - \text{Im}[S_f(u,v)] ,
\]

\[
S_f(u,v) = E_f(u,v) + jO_f(u,v) = E_h(u,v) - jO_h(u,v)
\]

\[
= \frac{1}{2} \left[ S_h(u,v) + S_h(-u,-v) \right].
\]

The generalisation of splitting single-variable functions into even and odd parts is the decomposition into symmetric and antisymmetric parts

\[
g(x,y) = g_{\text{sym}}(x,y) + g_{\text{antisym}}(x,y)
\]

with

\[
g_{\text{sym}}(x,y) = \frac{1}{2} \left[ g(x,y) + g(-x,-y) \right],
\]

\[
g_{\text{antisym}}(x,y) = \frac{1}{2} \left[ g(x,y) - g(-x,-y) \right].
\]

The real part of a 2D Fourier transform is symmetrical, and the imaginary part is antisymmetric. Therefore,

\[
E_f(u,v) = \frac{G_f(u,v) + G_h(-u,-v)}{2}
\]

\[
O_f(u,v) = \frac{G_f(u,v) - G_h(-u,-v)}{2}
\]

with \( G_f(u,v) = E_f(u,v) - jO_f(u,v) \), which is the 2D extension of equation (14).

The cross-correlation \( r(x,y) \) of the two functions \( p(x,y) \) and \( q(x,y) \) is

\[
r(x,y) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} p(x,y)q(x+x',y+y') \, dx \, dy ,
\]

where \( x' \) and \( y' \) are variable shift components along the \( x \)-direction and the \( y \)-direction respectively.

Evaluating the cross-correlation by means of the integral in equation (20) is complicated. This task can be significantly simplified in the Fourier space, where \( R_h(u,v) = P_h(u,v) \cdot Q_h(u,v) \); with \( R_h(u,v) \), \( P_h(u,v) \) and \( Q_h(u,v) \) being the Fourier transforms of \( r(x,y) \), \( p(x,y) \), and \( q(x,y) \), respectively. Therefore,

\[
r(x,y) = \text{FFT}^{-1}[R_h(u,v)] = \text{FFT}^{-1}[P_h^*(u,v) \cdot Q_h(u,v)] .
\]

In general, the presence of noise makes the 'exact' localisation of the cross-correlation peak difficult. In other words, using the POC function,

\[
r(x,y) = \text{FFT}^{-1} \left[ \frac{R_h(u,v)}{|R_h(u,v)|} \right]
\]

and we obtain a good compromise between peak sharpness and noise tolerance.

Using the Hartley transform, \( R_h \) can be written as [34]

\[
R_h(u,v) = \frac{R_h(u,v) + R_h(-u,-v) - iR_h(u,v) - R_h(-u,-v)}{2}
\]

\[
= \frac{P_h(u,v) + P_h(-u,-v) + iP_h(u,v) - P_h(-u,-v)}{2}
\]

\[
+ \frac{Q_h(u,v) + Q_h(-u,-v) - iQ_h(u,v) - Q_h(-u,-v)}{2}
\]

\[
= \frac{P_h(u,v) \cdot Q_h(u,v) + P_h(-u,-v) \cdot Q_h(-u,-v)}{2}
\]

where \( R_h(u,v) \) is the Hartley transform of \( r(x,y) \). From equation (23), we obtain

\[
R_h(u,v) + R_h(-u,-v)
\]

\[
= P_h(u,v) \cdot Q_h(u,v) - P_h(u,v) \cdot Q_h(-u,-v)
\]

\[
R_h(u,v) - R_h(-u,-v)
\]

\[
= P_h(u,v) \cdot Q_h(u,v) - P_h(-u,-v) \cdot Q_h(u,v)
\]

and eventually
\[ R_u(u,v) = \frac{P_u(u,v) \cdot Q_u(u,v) - P_u(-u,-v) \cdot Q_u(-u,-v)}{2} \]
\[ + \frac{P_u(-u,-v) \cdot Q_u(u,v) + P_u(-u,-v) \cdot Q_u(-u,-v)}{2} \] (25)

Therefore, the cross-correlation can be computed by means of the Hartley transform. Indicating with FHT, we have
\[ r(x,y) = \text{FHT}^{-1}\{R_u\} \]
\[ = \text{FHT}^{-1}\left\{ \frac{P_u(u,v) \cdot Q_u(u,v) - P_u(-u,-v) \cdot Q_u(-u,-v)}{2} \right\} \]
\[ + \frac{P_u(-u,-v) \cdot Q_u(u,v) + P_u(-u,-v) \cdot Q_u(-u,-v)}{2} \] (26)

Introducing the POC function, we have
\[ r(x,y) = \text{FHT}^{-1}\left\{ \frac{R_u(u,v)}{|R_u(u,v)|} \right\}. \] (27)

4. RESULTS AND DISCUSSION

A number of FHT algorithms have been developed [9], [35]-[39]. To test the proposed technique, an FHT was implemented by using the algorithms reported in [40]. Furthermore, the application of the Hartley POC function, in comparison with Fourier POC, was tested.

Figure 3 shows the POC obtained by two fingerprint patterns. In this example, the two images are each 512 x 512 pixels. The POC was computed in the Fourier space by means of equation (22) and in the Hartley space by means of equation (27). The FFT POC and the FHT POC are very similar, virtually identical.

In banknotes and other documents, metal fibres are present, as security features, inside the paper pulp. These fibres can only be observed under ultraviolet light. After image enhancement and some segmentation processes, it is possible to obtain a binary controlling grid that indicates the position of the majority security fibres. The distribution of these metallic fibres is random and not reproducible; each banknote can be told apart from the others using this random distribution. Obviously, the captured image has residual geometrical distortion and noise. Therefore, a verification approach based on digital phase correlation calculation can be a good solution [20], [21].

Figure 4 shows the POC used in the banknote identification. In this example, the fluorescent fibres (embedded in the paper) of two different 20-Euro banknotes are compared.

The POC functions obtained by means of FFT and FHT is practically identical. The examples illustrated do not prove that the FHT is superior to the FFT, but they do demonstrate that the FHT is fully compatible with the FFT.

Figure 5 shows an example of a drug package watermark realised with fluorescent points printed with ultraviolet light-sensitive ink. This watermark has a structure similar to the one used in speckle measurement [41]. In this example, the robustness of the POC approach against geometrical distortions and watermark partial disruption (also in the case of the Hartley transform POC) is highlighted. In particular, the correct watermark identification is highly relevant even in case c), where more than half of the watermark was unavailable.
In general, to perform FFT and FHT, the same number of operations are required. In other words, FHT requires a comparable number of steps to execute and is of comparable complexity to the FFT.

In any case, the FHT algorithm requires only real arithmetic computations compared to complex arithmetic operations in any standard FFT. Therefore, the speed of performing an FHT should be about twice as fast as the FFT. However, it is hard to compare the run times of these two programs for the transform process itself because the programs used in this effort are implemented differently.

On the other hand, the use of only real-value functions implies fewer computational resources and less computer memory; usually, the memory requirement by typical FFT algorithm is twice as high as for FHT.

5. CONCLUSION

In this paper, we have shown the correspondence between the POC function obtained by means of FFT and FHT.

The Hartley transform is similar to the Fourier transform, but it is free from the need to process complex numbers. Furthermore, the FHT performs the transformation itself, and the convolution and the cross-correlation are undertaken with fewer additions and multiplications than FFT. Using FHT to compute cross-correlation, there is no loss of information, and the results are identical to those obtained by FFT. Unfortunately, the numbers of all other operations are the same. Differences in time consuming exist between FHT and complex-valued FFT, and in any case, even compared with newly real-valued FFT algorithms, the possibility to use an identical programming solution for forward and inverse transformation is still an important point for using FHT instead of FFT in image correlation. In addition, the fact that the Hartley transform does not use complex numbers leads to the need for less memory to store numbers than the Fourier one. The Hartley transform can therefore be considered a valid alternative to FFT, particularly useful when a greater efficiency in memory requirements is needed, as, for instance, in embedded systems and smartphone apps.

REFERENCES


